

# Revisiting the microfield formulation using Continuous Time Random Walk Theory

D. Boland, A. Bouzaher, H. Capes, F. Catoire, M. Christova<sup>a</sup>, L. Godbert-Mouret, M. Koubiti, S. Mekkaoui, Y. Marandet, J. Rosato, R. Stamm

*Physique des Interactions Ioniques et Moléculaires  
Université de Provence et CNRS, Marseille, France*

<sup>a</sup>*Department of Applied Physics, Technical University-Sofia*



# Outline

**Microfield Model Method**

**The Continuous Time Random Walk Theory and its application to Stark Broadening**

**The Kangaroo Hypothesis**

**Time Memory effect**

**Application to Ly- $\alpha$**

**References**



# Semi-classical theory for line profile

$$i\hbar \frac{d}{dt} U(t, t') = (H_0 + V(t) + B)U(t, t')$$

$$V(t) = -\vec{D} \cdot \vec{E}(t)$$

$H_0$  – Hamiltonian operator for the unperturbed atom

$V(t)$  – time-dependent perturbing potential due to the motions of ions

$B$  – operator of electron impacts

## Assumptions:

- dipole approximation for the perturbing potential
- perturbers cause only radiative transitions
- perturbers are classical uncorrelated particles moving in straight lines

# Microfield Model Method

$$\frac{d}{dt} U(t, t') = M(\vec{E}(t)) U(t, t')$$

Linear stochastic equation of the evolution operator

$E(t)$  - fluctuating microfield

$$P(\vec{E})$$

$$\Gamma(t) = \langle \vec{E}(t) \cdot \vec{E}(0) \rangle$$

Probability distribution and covariance of the microfield → crucial role for Stark profiles

U. Frisch and A. Brissaud, J. Quant. Spectrosc. Radiat. Transfer 11, 1753 (1971)

C. Stehlé, Astron. Astrophys. Suppl. Ser. 104, 509 (1994)

B. Talin et al., Phys. Rev. A 51, 4917 (1995)

## This work

- To revise the **Stochastic Processes** used to model the plasma microfield, namely the **Kangaroo Process**, developed by *Frisch and Brissaud* (1971), at the light of the **Continuous Time Random Walk (CTRW)** of *Montroll, Weiss* (1965)
- Use the CTRW for calculating the ion dynamics effect on the Lyman alpha line
- Explore possible improvements of the stochastic process

E. Montroll and J. Weiss, J. of Mathematical Physics 6, 167 (1965)

# Continuous Time Random Walk (CTRW)



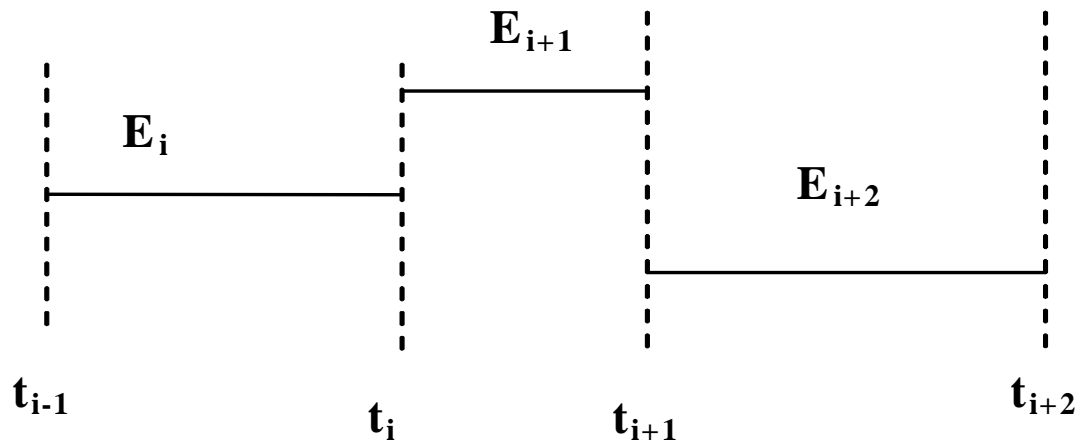
*Elliott Montroll  
(1916-1983)*

**CTRW**: the time interval between two jumps in the stochastic variable obeys a given probability law (Waiting Time Distribution)

**Kangaroo Process** is a particular case of CTRW process using the *Kangaroo Hypothesis* and a *Poisson* Waiting Time Distribution

Consider  $\mathbf{E}(t)$  as a stochastic vector

$$\mathbf{E}(t) \rightarrow \{\mathbf{E}_i, t_{i-1}\}$$



# Continuous Time Random Walk (CTRW)

*Key Quantity: probability density*

$$\Psi_{(t_i - t_{i-1})}(\mathbf{E}_{i+1}, \mathbf{E}_i)$$

Stationary process

Statistically independent jumps

$$\int d^3\mathbf{E} \int_0^{\infty} dt \psi_t(\mathbf{E}, \mathbf{E}') = 1$$

$$\int d^3\mathbf{E}' \psi_t(\mathbf{E}', \mathbf{E}) = \varphi_{\mathbf{E}}(t)$$

**Waiting Time Distribution**

$$\Phi_{\mathbf{E}}(t) = 1 - \int_0^t dt' \varphi_{\mathbf{E}}(t')$$

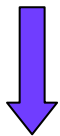
**Probability that no jumps occur during t**



# Solution of the Schrödinger equation for a stepwise constant process

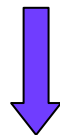
$$dU(t, t_0)/dt = -i(H_0 - D \cdot E(t))U(t, t_0)$$

$$U(t, t_0) = U_{E_{m+1}}(t - t_m) U_{E_m}(t_m - t_{m-1}) \dots U_{E_1}(t_1, t_0)$$



Single field history of an atom during t

$$U_{E_k}(t_k - t_{k-1}) = \exp[-i(H_0 - d \cdot E_k)(t_k - t_{k-1})]$$



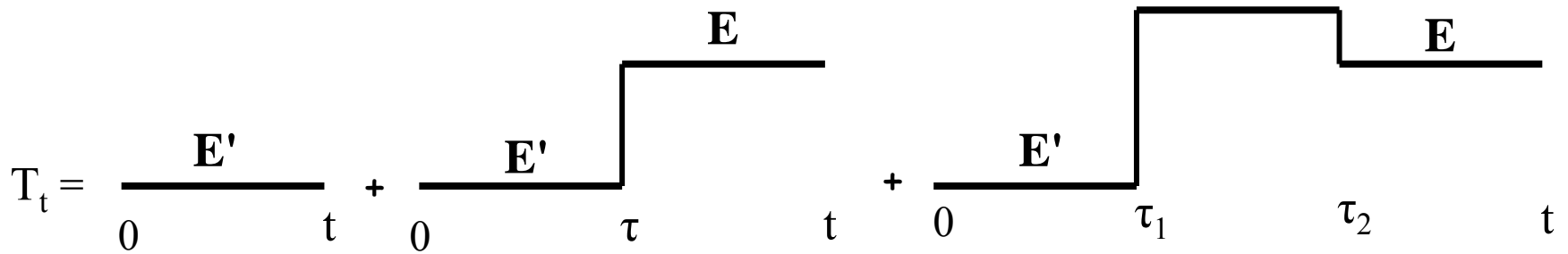
Time evolution for an atom in a static field  $E_k$





# Averaging of the evolution operator over electric fields

1) Calculate the atomic time evolution operator for all possible  $\mathbf{E}$  field time histories with the initial value  $\mathbf{E}'$  and final value  $\mathbf{E}$  at time  $t$   
 $T_t(\mathbf{E}, \mathbf{E}')$



$$T_t(\mathbf{E}, \mathbf{E}') = \delta(\mathbf{E} - \mathbf{E}') \Phi_{\mathbf{E}}(t) U_{\mathbf{E}}(t) + \int_0^t d\tau \int d\mathbf{E}_1 T_{t-\tau}(\mathbf{E}, \mathbf{E}_1) U_{\mathbf{E}'}(\tau) \Psi_{\tau}(\mathbf{E}_1, \mathbf{E}')$$

# Averaging of the evolution operator over electric fields

2) Average over all initial and final electric fields

Electric field Distribution

↓  
Hooper

$$T(t) = \int d^3E \int d^3E' P(E') T_t(E, E')$$



New part:

## Generalization of Kangaroo Process

Kangaroo Hypothesis



$$\psi_t(\mathbf{E}, \mathbf{E}') = q(\mathbf{E})\varphi_{\mathbf{E}'}(t)$$

$q(\mathbf{E})$  - the probability density to find an electric field  $\mathbf{E}$  after one jump

Waiting Time Distribution  $\varphi_{\mathbf{E}}(t)$  – any function

New part:

## Generalization of Kangaroo Process

**Exact solution for the atomic evolution operator**

$$\tilde{T}(s) = \int d^3E P(E) \tilde{Y}_E(s) + \int d^3E q(E) \tilde{Y}_E(s) \left[ I - \int d^3E q(E) \tilde{X}_E(s) \right]^{-1} \int d^3E P(E) \tilde{X}_E(s)$$

$$\tilde{X}_E(s) = \int_0^\infty dt \varphi_E(t) \exp - \left[ s + \frac{1}{i\hbar} (H_0 - \vec{D} \cdot \vec{E} + B) \right] t$$

$$\tilde{Y}_E(s) = \int_0^\infty dt \Phi_E(t) \exp - \left[ s + \frac{1}{i\hbar} (H_0 - \vec{D} \cdot \vec{E} + B) \right] t$$

# Standard Kangaroo process

$$\psi_t(\mathbf{E}, \mathbf{E}') = q(\mathbf{E}) \varphi_{\mathbf{E}'}(t)$$

$$\varphi_{\mathbf{E}}(t) = v(\mathbf{E}) \exp[-v(\mathbf{E})t]$$

In diagonal representation

$$p = \left[ s + \frac{1}{i\hbar} (H_0 - \vec{D} \cdot \vec{E} + B) \right]$$

$$\tilde{X}_{\mathbf{E}}(s) = \tilde{\varphi}_{\mathbf{E}}(p)$$

$$\tilde{Y}_{\mathbf{E}}(s) = \tilde{\Phi}_{\mathbf{E}}(p)$$

The Frisch-Brissaud's result for the atomic evolution operator  $T(t)$  is recovered if the Waiting Time Distribution  $\varphi_{\mathbf{E}}(t)$  is a Poisson's one

# Main difference from Frisch-Brissaud's solution

$$(\partial/\partial t + iH_0 + id.E)T_t(E, E_1) = -\int_0^t K_E(t-\tau)T_\tau(E, E_1)d\tau + q(E)\int_0^t d\tau \int d^3E' K_{E'}(t-\tau)T_\tau(E', E_1)$$

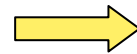
Memory kernel

$$\tilde{K}_E(s) = p\tilde{\varphi}_E(p)/(1-\tilde{\varphi}_E(p))$$

$$p = \left[ s + \frac{1}{i\hbar} (H_0 - \vec{D} \cdot \vec{E} + B) \right]$$

Poisson WTD

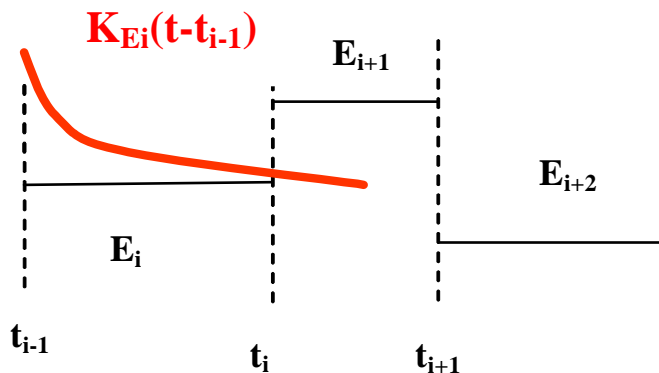
$$K_E(t) = v(E)\delta(t)$$



$$\tilde{K}_E(s) = v = \text{const}$$

Arbitrary WTD

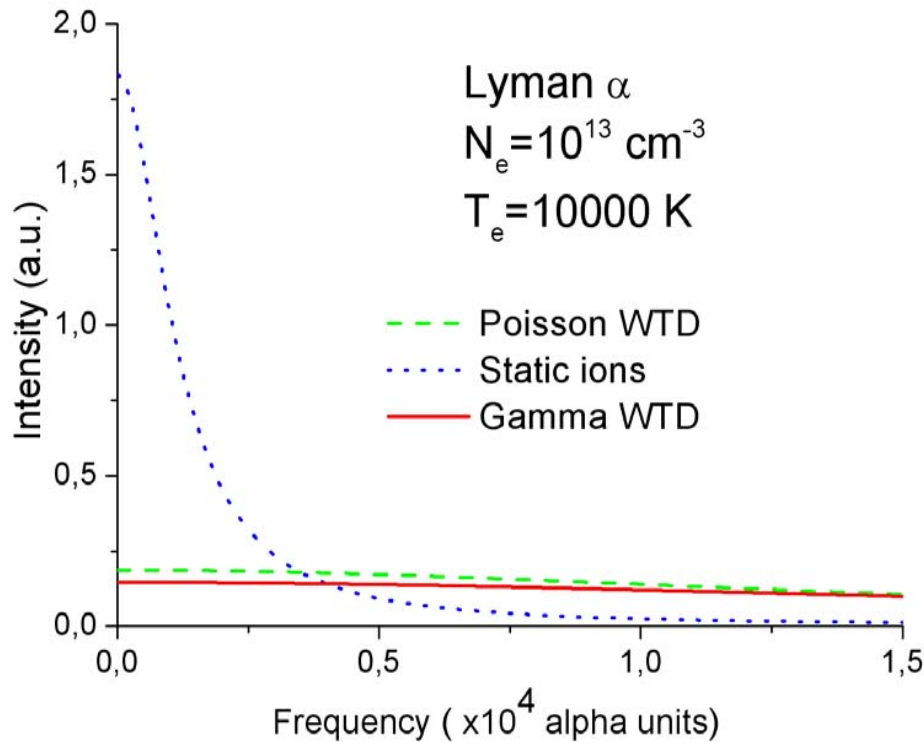
$K_E(t)$  non-local in time



jumps with different electric fields overlap in time

# Application to Ly-a line

Low density  
(near impact)



Effect of the electric field dynamics

$$q(E) = P(E)$$

$$\varphi_E(t) = \frac{v}{\Gamma(a)} (vt)^{a-1} \exp -vt$$

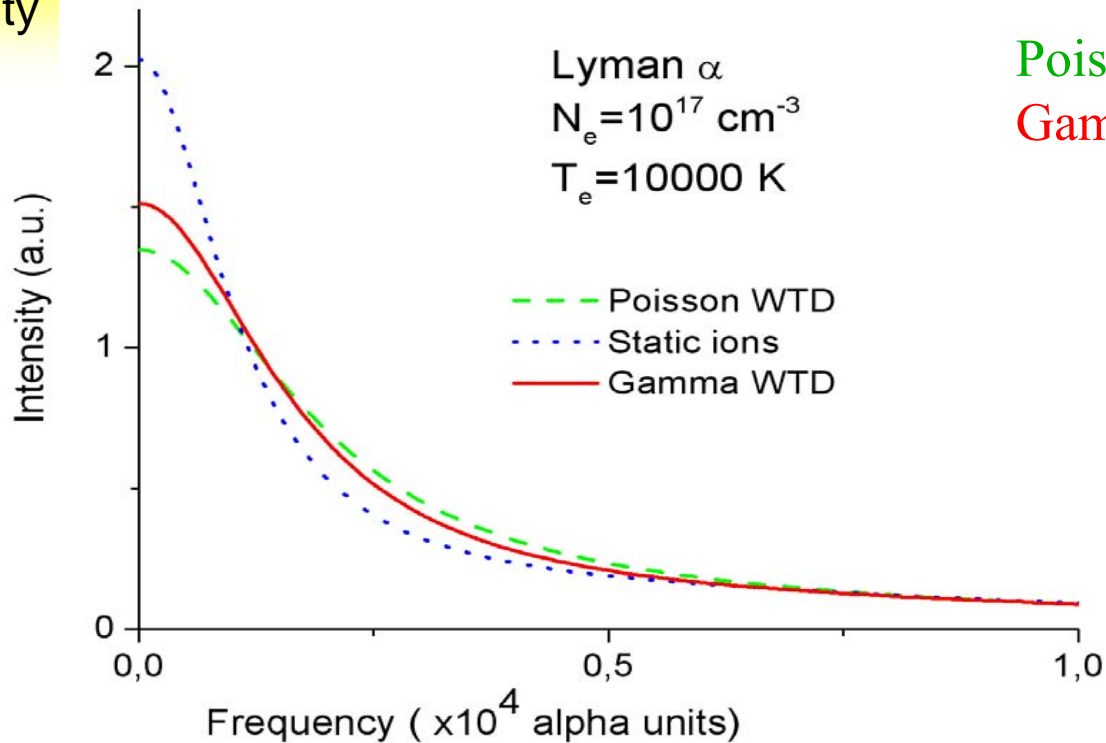
WTD

Poisson ( $a = 1$ )

Gamma ( $a = 2$ )

# Application to Ly-a line

High density  
(near static)



Poisson WTD  $a = 1$   
Gamma WTD  $a = 2$

Line profile with Gamma Waiting Time Distribution is close to the static profile



# Conclusions

- The Continuous Time Random Walk allows to generalize the Frisch Brissaud result to arbitrary Waiting Time Distribution
- The CTRW solution retains memory effects
- Preliminary applications to Ly- $\alpha$  show a significant influence of the choice of the Waiting Time Distribution

## Perspectives

### Stark Broadening

- Comparison with simulation results
- Use of Levy WTD for turbulent electric fields

Collisional Radiative Model with fluctuating plasmas parameters

Neutral Transport in turbulent plasmas

# References

- U. Frisch and A. Brissaud, J. Quant. Spectrosc. Radiat. Transfer 11, 1753 (1971)
- C. Stehlé, Astron. Astrophys. Suppl. Ser. 104, 509 (1994)
- B. Talin et al., Phys. Rev. A 51, 4917 (1995).
- E. Montroll and J. Weiss, J. of Mathematical Physics 6, 167 (1965)

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